



On some Gibbs measures of a coupled Ising-Ising model

Dadamirzayeva I.X

Namangan State University, Namangan, Uzbekistan

iqboladadamirzayeva718@gmail.com

Abstract: In this paper, we study the translation-invariant Gibbs measures of a coupled Ising-Ising model defined on a Cayley tree of order $k \geq 2$. The model is derived as a special case of the coupled Ising-Potts model when the number of Potts states $q = 2$. We consider a Hamiltonian with spin configurations represented by $(s, \sigma) \in \{-1, 1\} \times \{1, 2\}$ and obtain a system of nonlinear equations characterizing the corresponding Gibbs measures. For the particular case $k = 2$, the system is analyzed in detail. We prove that if the interaction parameter satisfies the condition $\theta + \theta^{-1} > 2 + 2$, then the system admits at least three distinct translation-invariant Gibbs measures. This result highlights the occurrence of a phase transition in the model and contributes to the understanding of multi-component spin systems on hierarchical structures.

Keywords: Gibbs measures, coupled Ising-Ising model, Cayley tree, translation-invariance, Hamiltonian, phase transition, statistical mechanics.

Here we give a result related to Gibbs measures of a Hamiltonian defined on a Cayley tree $\Gamma^k = (V, L)$ (see [1] for main definitions). Namely, following [2] we consider a coupled Ising-Potts model on Cayley trees of order $k \geq 2$. This has spin vectors (s, σ) , where $s \in I = \{-1, 1\}$ and $\sigma \in \Phi = \{1, \dots, q\}$, $q \geq 2$.

A coupled configuration (s, σ) on V is then defined as a function $x \in V \rightarrow (s(x), \sigma(x)) \in I \times \Phi$. The following Hamiltonian is called the coupled Ising-Potts model:



$$H(s, \sigma) = -J \sum_{\langle x, y \rangle \in E} s(x)s(y)\delta_{\sigma(x)\sigma(y)} \quad (1)$$

where $J \in \mathbb{R}$ is a constant, $\langle x, y \rangle$ is nearest neighbor vertices and δ_{ij} is the Kroneker’s symbol.

In [2] it is proven that to each Gibbs measure of Hamiltonian corresponds a vector-valued function $h_x = (h_{\epsilon, i, x})_{(\epsilon, i) \in I \times \Phi}$ defined on vertices of the Cayley tree that satisfies for any $x \in V$ the following equations:

$$z_{\epsilon, i, x} = \prod_{y \in S(x)} \frac{(1 - \theta^{-\epsilon})(\theta^{\epsilon} z_{1, i, y} - z_{-1, i, y}) + \sum_{j=1}^{q-1} (z_{-1, j, y} + z_{1, j, y}) + z_{1, q, y} + 1}{\theta + \theta^{-1} z_{1, q, y} + \sum_{j=1}^{q-1} (z_{-1, j, y} + z_{1, j, y})}, \quad (2)$$

$$z_{1, q, x} = \prod_{y \in S(x)} \frac{(1 - \theta_{-1})(\theta z_{1, q, y} - 1) + \sum_{j=1}^{q-1} (z_{-1, j, y} + z_{1, j, y}) + z_{1, q, y} + 1}{\theta + \theta^{-1} z_{1, q, y} + \sum_{j=1}^{q-1} (z_{-1, j, y} + z_{1, j, y})}, \quad (3)$$

where $\epsilon \in I$, $z_{\epsilon, i, x} = \exp\{h_{\epsilon, i, x} - h_{-1, q, x}\}$, $i = 1, \dots, q - 1$, $\theta = \exp(J\beta)$, and $S(x)$ is the set of direct successors of x .

We consider Gibbs measures which are translation-invariant, i.e correspond to solutions of the form:

$$z_{\epsilon, i, x} \equiv z_{\epsilon, i}, \text{ for all } x \in V.$$

Here we consider the case $q = 2$, $k = 2$. Note that in case $q = 2$ the Potts model coincides with the Ising model. This is why we call our model as an Ising-Ising model. Then by (2), (3) denoting $u_1 = z_{-1, 1}$, $v_1 = z_{1, 1}$, $v_2 = z_{1, 2}$ we get:

$$u_1 = \left(\frac{(1 - \theta^{-1})(\theta u_1 - v_1) + u_1 + v_1 + v_2 + 1}{\theta + \theta^{-1} v_2 + u_1 + v_1} \right)^2, \quad (5)$$

$$v_1 = \left(\frac{(1 - \theta^{-1})(\theta v_1 - u_1) + u_1 + v_1 + v_2 + 1}{\theta + \theta^{-1} v_2 + u_1 + v_1} \right)^2, \quad (6)$$



$$v_2 = \left(\frac{(1 - \theta^{-1})(\theta v_2 - 1) + u_1 + v_1 + v_2 + 1}{\theta + \theta^{-1}v_2 + u_1 + v_1} \right)^2. \quad (7)$$

Through making equation simpler we get:

$$u_1 = \left(\frac{\theta^{-1}v_1 + \theta u_1 + v_2 + 1}{\theta + \theta^{-1}v_2 + u_1 + v_1} \right)^2, \quad (8)$$

$$v_1 = \left(\frac{\theta^{-1}u_1 + \theta v_1 + v_2 + 1}{\theta + \theta^{-1}v_2 + u_1 + v_1} \right)^2, \quad (9)$$

$$v_2 = \left(\frac{\theta^{-1} + \theta v_2 + u_1 + v_1}{\theta + \theta^{-1}v_2 + u_1 + v_1} \right)^2. \quad (10)$$

Let $u_1 = u$, $v_1 = v$, and $v_2 = 1$. We have following system of equations:

$$\begin{cases} u = \left(\frac{\theta u + \theta^{-1}v + 2}{\theta + \theta^{-1} + u + v} \right)^2 \\ v = \left(\frac{\theta^{-1}u + \theta v + 2}{\theta + \theta^{-1} + u + v} \right)^2 \end{cases};$$

In the case when $u = v$, we have to solve the following equation from the system.

$$u = \left(\frac{\theta^{-1}u + \theta u + 2}{\theta + \theta^{-1} + 2u} \right)^2 \quad (11)$$

Let's denote $\theta + \theta^{-1}$ as a and take the square root of both sides of the equation. As a result, we will obtain the following equation:

$$\sqrt{u} = \frac{au + 2}{a + 2u} \quad (12)$$

Let $\sqrt{u} = x$ and we have following equation

$$2x^3 + ax - ax^2 - 2 = 0 \quad (13)$$

$$2(x^3 - 1) + a(x - x^2) = 0 \quad (14)$$

$$(x - 1)(2x^2 + 2x + 1 - ax) = 0$$



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As a result we get three $x_1 = 1$, $x_2 = \frac{a-2-\sqrt{a^2-4a-4}}{4}$ solutions

$$x_3 = \frac{a-2+\sqrt{a^2-4a-4}}{4} \quad \text{and}$$

Considering that $u = x^2$ analyzing the expression under the square root is enough. $a > 2 + 2$, i.e $\theta + \theta^{-1} > 2 + 2$, system has three solutions: (1,1,1)

$$\left(\left(\frac{a-2-\sqrt{a^2-4a-4}}{4} \right)^2, \left(\frac{a-2-\sqrt{a^2-4a-4}}{4} \right)^2, 1 \right) \quad \text{and}$$
$$\left(\left(\frac{a-2+\sqrt{a^2-4a-4}}{4} \right)^2, \left(\frac{a-2+\sqrt{a^2-4a-4}}{4} \right)^2, 1 \right)$$

As mentioned above to each solution of the system (2), (3) corresponds a Gibbs measure.

Therefore we have the following result:

Theorem 1. For the coupled Ising-Ising model, if $\theta + \theta^{-1} > 2 + \sqrt{2}$ then there are at least three translation-invariant Gibbs measures.

References:

1. Rozikov U.A, *Gibbs measures in biology and physics: The Potts model.*, World Scientific, Singapore, 2023.
2. Haydarov F.H, Omirov B.A, Rozikov U.A, *Coupled Ising-Potts Model: Rich set of critical temperatures and translation-invariant Gibbs measures.*, arXiv:2502.12014v11. 2025.